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Citation: Applied Physics Letters 101, 193902 (2012); doi: 10.1063/1.4766736
View online: http://dx.doi.org/10.1063/1.4766736
View Table of Contents: http://scitation.aip.org/content/aip/journal/apl/101/19?ver=pdfcov
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Plasmonic-enhanced performance for polymer solar cells prepared with inverted structures

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(Received 21 September 2012; accepted 25 October 2012; published online 9 November 2012)

We incorporated gold nanoparticles (Au NPs) in inverted organic photovoltaic devices to enhance the device performance. The photocurrent and fill factors were improved after the addition of Au NPs into the Cs2CO3 buffer layer. The photoluminescent measurements revealed a significant increase of light absorption of the photoactive layer. We attribute the improvement to local field enhancement induced by the localized surface plasmon resonance. Further, through the study of the morphologies of the cathode interfaces, we found that the rough surfaces might increase the device resistances. This drawback, however, was overwhelmed by the advantageous plasmonic effects. © 2012 American Institute of Physics.

Organic photovoltaic devices (OPVs) have been considered as one of the most promising systems for harnessing solar energy because of their light weight, mechanical flexibility, and the ability to prepare large-area panels at low cost.1,2 At present, OPVs based on the concept of bulk-heterogeneous junction can exhibit power conversion efficiencies (PCE) as high as 10%.3 The external quantum efficiency of OPVs is governed by its internal quantum efficiency and absorption efficiency.4,5 Because the former, which is affected by the diffusion and dissociation of excitons and charge collection, can now approach almost 100%,3 the absorption efficiency remains one of the major problems toward even higher efficiencies.6 One feasible approach for increasing the absorption efficiency is to use a thicker photoactive layer. Nevertheless, due to the low carrier mobilities of organic materials, the use of too thick films inevitably increases the possibility of charge recombination and device resistance, thereby decreasing the charge collection efficiency.6–8 As a result, the challenge remains to achieve OPVs with high absorption efficiencies without degrading their internal quantum efficiencies.

Light trapping based on surface plasmon resonance (SPR) effects has attracted much attention as a means of improving the efficiency of OPVs.7–15 Various nanostructures, such as metallic nanoparticles, have been introduced into the devices for triggering SPRs.7–15 Their unique optical properties, including local field enhancement and strong light scattering, might improve absorption processes in OPVs. As for the device structures, most plasmonic-enhanced OPVs reported to date have been conventional ones. Because inverted device architectures eliminated the use of low-work-function metals, which are air-sensitive, they usually exhibited prolonged device lifetimes.16–18 Plasmonic-enhanced OPVs possessing an inverted structure, however, are still very rare. In this work, we incorporated gold nanoparticles (Au NPs) into inverted OPVs to increase the light absorption efficiency. The Au NPs induced localized surface plasmon resonance (LSPR), which enhanced the local electromagnetic field of the active layer, thereby improving the device performance.

The devices were fabricated on indium tin oxide (ITO)-coated substrates. The device structure in this study is illustrated in Fig. 1(a). The ITO glasses were treated with UV-ozone prior to use. The Au NP solution was prepared

![Fig. 1. Device structures of the OPVs in this study. (b) SEM image of the Au NPs. (c) J-V characteristics, recorded under illumination at 100 mW cm−2 (AM 1.5G), of the OPVs prepared with Cs2CO3 layers incorporating various amounts of Au NPs.](http://dx.doi.org/10.1063/1.4766736)
using procedures described previously. Fig. 1(b) displays the scanning electron microscopy (SEM) image of the Au NPs. The average particle size was ca. 45 nm. To prepare the cathodic buffer layer, the NP solution was blended into the Cs2CO3 solution with various volume ratios. Then, the composite solution was spin-coated onto the substrates, and the resulting thin film was annealed at 140 °C for 15 min. The photocative layer comprised regioregular poly(3-hexylthiophene) (P3HT) and 1-(3-methoxycarbonyl)propyl-1-phenyl[6,6]methanofullerene (PCBM) at a weight ratio of 1:1. The photocative film was spin-coated from a solution of 1,2-dichlorobenzene on the Cs2CO3 layer. After solvent annealing, the dried film was further annealed at 110 °C for 15 min. The thickness of the active layer was ~180 nm. Finally, a bilayer anode comprising MoO3 (5 nm) and Ag (100 nm) was deposited through thermal evaporation. For the device characteristics, the photocurrent density-voltage (J-V) curves under illumination were measured using a Keithley 2400 source measure unit. The light source was a 150 W Thermal Oriel solar simulator. The IPCE measurements system (Enli Technology) comprised a quartz-tungsten-halogen lamp, a monochromator, an optical chopper, a lock-in amplifier, and a calibrated silicon-based diode.

Figure 1(c) displays the J-V characteristics of the OPVs prepared with various amounts of Au NPs. The device prepared with pristine Cs2CO3 exhibited an open-circuit voltage (Voc) of 0.55 V, a short-circuit current (Jsc) of 9.73 mA cm−2, and a fill factor (FF) of 0.58, resulting in a PCE of 3.12%. After the addition of Au NPs to the Cs2CO3 layer, the values of Voc remained unchanged (0.55 V), suggesting that the cathodic interface was not affected too much by the Au NPs. Some portion of the surface of the NPs was probably also modified by Cs2CO3. The photocurrent and FFs, however, were improved (Table I). For the devices prepared with 20% Au NP solution, the Jsc and FF values were 10.11 mA cm−2 and 0.64, respectively. As a result, the PCE was increased to 3.54%. Nevertheless, a further increase in the amount of the Au NP in the Cs2CO3 layer resulted in a decrease in the value of Jsc, presumably due to enhanced backward scattering and/or increased resistivity of the buffer layer. Notably, the values of the device series resistance (Rs), extracted from the J-V curves in the dark, indeed increased with the increasing Au NP concentrations. The results of extraction of Rs, listed in Table I, reveal that the enhanced device performance did not result from a reduction in device resistance.

Figure 2 displays the incident photon-to-electron conversion efficiency (IPCE) curves for various devices. Clearly, the photocurrent was increased after incorporating the Au NPs. One can also see that the trends in the IPCE follow those for the values of Jsc. We further compared the curve of the increase in IPCE (ΔIPCE) after the addition of 20% Au NPs with the extinction spectrum of the NPs (Fig. 2(b)). The wavelength regime, in which the IPCE values were increased, coincides with the extinction range of Au NPs, suggesting that excitation of LSPR indeed improved the efficiencies. Note that the extinction spectrum of the dilute Au NPs was relatively narrow. The real resonance peak of NPs should strongly depend on their surrounding media.

To understand the reason for the increased value of Rs of the device prepared with a higher concentration of Au NPs, we used atomic force microscopy (AFM) to observe the surface morphologies of the Cs2CO3 layers. Fig. 3 displays the AFM images of the buffer layers containing various amounts of Au NPs. We observed significant morphological changes on the surfaces after embedding the Au NPs. Further, the root-mean-square (rms) roughness of the surface prepared without Au NPs was 2.29 nm; the roughness increased with the concentration of an NP solution. They were 4.30 nm, 4.78 nm, and 5.23 nm for the samples fabricated with 10%, 20%, and 30% of NP solutions. These results suggest that the rough surfaces might induce trap states at the contacts and increase the device resistances. Fortunately, this drawback was overwhelmed by the advantageous plasmonic effects, thereby improving the overall device PCEs.
Further, we performed the steady state photoluminescence (PL) measurements of the P3HT:PCBM thin films prepared on various Cs$_2$CO$_3$ buffer layers. The peak fluorescence intensity of the plasmonic sample was apparently enhanced by ca. 38% (Fig. 4). Because the resonance peak of the Au NPs was close to the absorption band of P3HT, we inferred that the enhanced PL was probably due to the increased level of photon absorption. Local enhancement of the electromagnetic field surrounding the Au NPs helped the energy dissipation, thereby increasing the total number of excitons in the photoactive layer.8

Previously, we blended Au NPs into the poly(3,4-ethylenedioxythiophene):polystyrenesulfonate (PEDOT:PSS) buffer layer to improve the device efficiencies.7,8 We attributed the improvement in device performance to the plasmonic effects. Another possible factor responsible for the increased photocurrent was related to the increased conductivity of the PEDOT:PSS layer after the addition of Au NPs.24 The trace surfactants left behind the synthesis process of the Au NP solutions could change the morphology of PEDOT:PSS and increase its conductivity.24 Nevertheless, this approach presented herein for constructing plasmonic-enhanced OPVs did not involve in the use of PEDOT:PSS. From the calculated series resistances, the cathode contacts were even deteriorated after the addition of Au NPs (Table I). Therefore, the results of this study can exclude the influence of the PEDOT:PSS layer and further confirm the plasmonic-enhanced mechanism of the Au NPs on the device performance.

In summary, we have improved the efficiency of inverted OPVs by incorporating Au NPs into the cathode buffer layer. The results of PL measurements revealed a significant increase of light absorption of the photoactive layer. The primary origin of the device improvement was local field enhancement induced by the LSPR. Through the study of the morphologies of the cathode interfaces, we found that the rough surfaces might increase the device resistances. Fortunately, this drawback was overwhelmed by the advantageous plasmonic effects. Because inverted devices usually exhibit longer device lifetime, we foresee that the results reported herein might be of further use in other material systems, such as low-band-gap polymers, to achieve even higher PCEs and enhanced stability.
We thank the National Science Council of Taiwan (Grant Nos. NSC 101-3113-E-009-005 and NSC 101-2628-E-009-008-MY3) and the Ministry of Education of Taiwan (through the ATU program) for financial support.