Improved Near-infrared Luminescence of Si-rich SiO$_2$ with Buried Si Nanocrystals Grown by PECVD at Optimized N$_2$O Fluence

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ABSTRACT

The optimized N$_2$O fluence for plasma enhanced chemical vapor deposition (PECVD) growing silicon-rich substoichiometric silicon oxide (SiO$_x$) with buried Si nanocrystals is demonstrated. Strong room-temperature photoluminescence (PL) at 550-870 nm has been observed in SiO$_x$ thin films grown by PECVD with N$_2$O fluence varying from 105 to 130 sccm. After annealing from 15 to 180 min, a 22-nm-redshift of the PL has been detected. The maximum PL intensity is observed for the 30-min annealed SiO$_x$ growing at N$_2$O fluence at 120 sccm. Larger N$_2$O fluence and longer annealing time causes a PL blueshift by 65 nm and 20 nm, respectively. Such a blueshift is attributed to shrinkage in the size of the Si nanocrystals under the participation of dissolved oxygen atoms from N$_2$O. The (220)-oriented Si nanocrystals with radius ranging from 4.4 to 5.0 nm are determined. The luminescent lifetimes lengthens from 20 $\mu$s to 52 $\mu$s as the nc-Si size extends from 4.0 to 4.2 nm. Optimal annealing times for SiO$_x$ preparing at different N$_2$O fluences and an optimum N$_2$O fluence of 120 sccm are reported. Serious oxidation effect at larger N$_2$O fluence condition is observed, providing smaller PL intensity at shorter wavelengths. In contrast, the larger size nc-Si will be precipitated when N$_2$O fluence becomes smaller, leading to a weaker PL at longer wavelength. These results provide the optimized growth condition for the Si-rich SiO$_2$ with buried Si nanocrystals.

Keywords: Si-rich SiO$_2$, photoluminescence, Nanocrystallite silicon, TEM, PECVD, TRPL

1. INTRODUCTION

Silicon nanocrystals (nc-Si) have been precipitated from the Si-rich silicon dioxide (SRSO) layer which is produced by several techniques such as high-dosage Si-ion-implantation in SiO$_2$, laser ablation, electron-beam evaporation, sputtering and plasma enhanced chemical vapor deposition (PECVD). It has been demonstrated that the photoluminescence (PL) from the nc-Si was enhanced with the increasing the annealing-temperature. Moreover, the peak wavelength slightly redshifts with the increasing temperature of the thermal treatment. Such results have been explained by quantum confinement theory. The quantum-confined nc-Si enlarge the band gap and increase the spatial overlaps of the electron and hole wave functions in Si, thus increasing the spontaneous emission rate and shifting the emission peak to energies higher than the crystalline Si band gap. Iacona et al. demonstrated that the PL wavelength is increasing from 770 nm to 880 nm as the annealing temperature rises from 1100°C to 1300°C. They also demonstrated the change of $\lambda_{\text{max}}$ (from 740 nm to 910 nm) at varied N$_2$O/SiH$_4$ flow ratio (ranging between 6 and 15), which correlates to the increase in Si composition in SiO$_x$ films between 33 and 44 at. %. The PL intensity of the 1200°C-annealed SiO$_x$ sample with 42 at. % of Si content is highest. Oliveira et al. also observed that demonstrated by varying the N$_2$O/SiH$_4$ flow ratio between 0.25 and 0.5 and keeping SiH$_4$ flow of 15 sccm. After deposition, the films were thermally annealed in N$_2$ ambient, at temperatures from 450 to 1000°C for 2 h. The PL spectra for the as-deposited samples exhibit a broad emission band in 1.5-2.1 eV range. The contribution of two peaks can be distinguished at $\sim$ 1.6 and 1.9 eV. After heat treatment, the PL of the samples grown with R = 0.3 and 0.5 were
observed that the overall emission is more intense in the as-deposited films for all the samples studied. For the sample grown with $R = 0.3$, which has a higher Si content, both PL peaks decrease and shift towards lower energy after the annealing steps at 450 and 550°C. After the heat treatment at 750°C, the lower energy peak vanishes, the one at $\sim 1.8$ eV increases and a lower intensity peak around 2.1 eV appears. After the annealing at 1000°C, the PL intensity decreases and tends to energize higher than 2 eV.

Varying the N$_2$/SiH$_4$ flow ratio is known to control the Si content, however, the correlation between N$_2$/SiH$_4$ ratio and excess Si density as well as nanocrystallite Si size is not clear. In particular, the optimized annealing condition could also be changed for the SiO$_x$ grown different N$_2$/SiH$_4$ ratios. Since the control of SiH$_4$ fluence is unavailable to give a predictable result according to previous results, we investigate in this work the effect of N$_2$/O fluence and annealing time on the PL intensity and wavelength. We also investigate the visible PL and time-resolved photoluminescence (TRPL) from PECVD SiO$_x$ thin films after post-deposition annealing in N$_2$ ambient by varying the annealing times under the optimal N$_2$/O fluence. The morphology of PECVD-grown SiO$_x$ thin film was investigated using transmission electron microscope (TEM) to support the existence of nc-Si.

2. EXPERIMENT

The SiO$_x$ films were grown on p-type Si (100) substrate by using PECVD system at pressure of 70 mTorr with different SiH$_4$ and N$_2$/O fluences under forward power of 60 W. The N$_2$/O fluence was varied from 105 sccm to 130 sccm. After deposition, the samples were annealed in a quartz furnace at 1100°C under N$_2$ flowing for 15 min to 180 min. The structure and optical properties of these films have been studied by TEM, PL and time-resolved photoluminescence (TRPL). The room temperature PL of the SiO$_x$ films used a He-Cd laser at wavelength and average power intensity of 325 nm and 5 W/cm$^2$ and analyzed with a fluorescence spectrophotometer (Jobin Yvon, TRIXA-320 with resolution of 0.06 nm) and a cooled photomultiplier (Jobin Yvon, Model 1424M). In order to support the existence of nc-Si, the bright-field cross-section viewing photograph had been taken by using TEM (JEOL 4000EX) operating at 400 keV with point-to-point resolution of 0.18 nm. In the TRPL experiment, the samples were pumped by a YAG laser (NY 60, Countinous) at 355 nm (the repetition rate is 1 Hz, the pulse full width of half maximum (FWHM) is 60 ps, and average power is 0.5 mJ/pulse) and detected by a time-correlated single-photon counting system at wavelength corresponding to the PL of nc-Si ($\lambda = 760$ nm).

According to Einstein’s two-level quantized radiation model, the PL intensity of nc-Si can be approximated by $I = \eta \sigma G N / \tau$, where $\sigma$ is the absorption cross-section of nc-Si that can be calculated by a theoretical approach of $\sigma = \tilde{\varepsilon} / (8\pi \nu \Delta \nu)$, $\eta$ is a relative coefficient, $\tau$ is the lifetime of nc-Si, $G$ is the pumping flux and $N$ is the nc-Si concentration$^{12-14}$. In the TRPL experiment, the correlation among the emission (absorption) cross-section, density and lifetime of nc-Si has to be obtained from the rate equation of a two-level system.

$$\frac{dN_1}{dt} = \frac{dN_2}{dt} = -A_{21}N_2 + B_{12}N_1 \rho(\nu) - B_{21}N_2 \rho(\nu), \quad (1)$$

where $N_1$ and $N_2$ are the population density in the state 1 and state 2, $A_{21}$ is the spontaneous emission coefficient, $B_{21}$ is the stimulated emission coefficient and $\rho(\nu)$ is the radiation density per frequency interval.

Since there is no stimulated emission in the TRPL analysis, the rate equation is simplified as

$$\frac{dN_1}{dt} = \frac{dN_2}{dt} = -A_{21}N_2 + B_{12}N_1 \rho(\nu), \quad (2)$$

Assuming that $N_1$ equals to nc-Si and $A_{21}$ equals to a reciprocal lifetime, in the steady state, the electron density in the state 2 ($N_2$),

$$N_2 = \frac{B_{12}}{A_{21}} N_1 \rho(\nu) = B_{12} \tau_{nc-Si} \rho(\nu) N_{nc-Si}, \quad (3)$$

According to the Einstein relations of $B_{21}=B_{12}$ and

$$B_{21} = \frac{c^3 A_{21}}{8 \pi n^+ n \hbar \nu}, \quad (4)$$
By expressing \( \rho(v) = \int \phi(t) dt \ A \) (where \( A \) is the spot size of the pumping laser), we have

\[
I_{PL} = \frac{c^3 A_{21}^2}{8\pi n^2 n_g \, \hbar^3} \, \rho(v) \int dt \ A \tau_{nc-Si} \, N_{nc-Si},
\]  

(4)

Assuming a rectangular PL function of \( I_{PL} \) and a rectangular pump function of \( \phi(t) \), we can rewrite Eq. (5) as

\[
I_{PL} \tau_{PL} \approx \frac{c^3}{8\pi n^2 n_g \, \hbar^3} \, \phi(t) \, N_{nc-Si} \, \tau_{pump},
\]  

(6)

\[
I_{PL} = \eta \, \sigma \, N_{nc-Si} \, \phi(t) \frac{\tau_{pump}}{\tau_{PL}} \propto \sigma \phi(t) \tau_{PL}^{-1} N_{nc-Si},
\]  

(7)

The theoretically equation (7) is also obtained in other papers\textsuperscript{14-16}, as expressed by \( I \propto \sigma \phi(t) \int \frac{\tau_{PL}}{\tau_R} \), for estimating the variations of nc-Si density through the PL intensity and lifetime.

### 3. RESULTS AND DISCUSSION

The PL spectra of PECVD-grown SiO\textsubscript{x} films under 120-sccm N\textsubscript{2}O-fluence as a function of annealing time were shown in Fig. 1. It was found that there is a broadened PL spectrum between 400 nm and 650 nm, which is attributed to the radiative defects including neutral oxygen vacancy (NOV)\textsuperscript{17} and the non-bridging oxygen hole center (NBOHC)\textsuperscript{18} in the as-grown SiO\textsubscript{x} sample.

![PL spectra](image)

Fig. 1. Annealing-time dependent of PL intensities under the N\textsubscript{2}O fluence at 120 sccm.

After annealing at 1100°C for 15 min, the defect-related PL is eliminated and the PL at 735 nm was enhanced. The central wavelengths and their FWHM of PL spectra for samples annealing from 15 min to 180 min were changed from 735 nm to 761 nm and 132.5 nm to 177 nm, respectively. In particular, the strongest PL appears in the 30-min annealed sample. The red shift in PL for the 30-min annealed sample was about 26 nm; however, a blue-shift phenomenon is observed after longer annealing time. In a 180-min annealed sample, the blue shift is up to 16.5 nm. Previously, Iacona \textit{et al.}\textsuperscript{8} also found the strong PL in the wavelength ranging from 650 nm to 950 nm, which is attributed to the nc-Si clusters in high-temperature annealed PECVD-grown SiO\textsubscript{x} film.
The nc-Si diameter as a function of annealing time by the rule of \( \frac{\lambda}{d} = \frac{1.24}{(1.12 + 3.73/d^{0.5})} \) was shown in Fig. 2. From the above equation, the size of nc-Si is proportional directly to the central wavelength of PL spectra. It shows that the thermal annealing process helps the formation and the size increasing of nc-Si during the first annealing process, so the red shift in PL was shown. This leads to the increasing quantity and wavelength of nc-Si with lengthening annealing time during the first 30 min. Later on, the oxygen atoms bond with the silicon atoms of nc-Si, which not only cause the reduction in nc-Si size, but also decrease the nc-Si density. As a result, the PL wavelength blue shifts, its FWHM enlarges and the PL intensity becomes weaker slowly.

In the inchoative experiment, the SiO\(_x\) samples were manufactured by changing one of the SiH\(_4\) and N\(_2\)O fluences. However, the PL spectrum was unavailable to give a predictable result for samples with varying SiH\(_4\) fluences. Therefore, the SiO\(_x\) samples were prepared by changing N\(_2\)O fluences from 105 sccm to 130 sccm, and the thick were in the range from 195 nm to 258 nm. The SiO\(_x\) samples were annealed at 1100°C for 30 min to obtain a maximum PL intensity, as shown in figure 3. As the N\(_2\)O fluence increased from 105 sccm to 120 sccm, the PL intensity at 120 sccm is enhanced 4 times larger than that at 105 sccm, but the PL intensity is decreased 3 times while the N\(_2\)O fluence increased to 130 sccm. The central PL wavelength decreases to 606.5-624.5 nm when the N\(_2\)O fluence is higher or lower than the optimum fluence. As the N\(_2\)O fluence detunes larger than the optimized N\(_2\)O fluence, the oxidation of excess silicon atoms become significant, which degrade the nc-Si related PL signal. Both the size and PL intensity of nc-Si become smaller in this case. In contrast, more excess Si atoms will exhibit in the SiO\(_x\) film as the N\(_2\)O fluence less than the optimal one, which leads to an increasing size of the buried nc-Si, and the larger size of nc-Si causes a smaller density as well as a smaller PL intensity. Nonetheless, the PL wavelength of the N\(_2\)O fluence less than the optimal one was not larger than the N\(_2\)O fluence more than the optimal one. It shows the correlative size of nc-Si was smaller than the latter one (the N\(_2\)O fluence more than 120 sccm).
In Fig. 1, it is evident that a sufficient annealing time to precipitate the nc-Si with suitable size and maximum density. Note that the annealing time may be varied for samples growing with different N\textsubscript{2}O fluence. This is the reason between the variance of the correlative sizes as the N\textsubscript{2}O fluences less or more than 120 sccm. The optimized annealing times of samples preparing under different N\textsubscript{2}O fluences were shown in Fig. 4.

The annealing time become longer (60-210 min) for smaller N\textsubscript{2}O-fluence conditions, whereas it is shorter than 30 min for the N\textsubscript{2}O fluence >120 sccm. The PL intensity and wavelength of all samples at the optimal annealing time under different N\textsubscript{2}O fluences are shown in Fig. 5. The PL wavelength first decreases from 733 nm to 754.5 nm as the N\textsubscript{2}O fluence increases from 105 sccm to 120 sccm, and then decreases to 690 nm at N\textsubscript{2}O fluence of 130 sccm or larger. The size of nc-Si becomes larger as the N\textsubscript{2}O fluence is smaller than optimum value (120 sccm), which is attributed to the more excess Si under less N\textsubscript{2}O fluence. The evolutions of PL intensity and wavelength under different N\textsubscript{2}O fluences in Fig. 5 were correlated to the above discussion.
The bright-field cross-section viewing photograph of TEM analysis for the optimized sample (the N₂O fluence at 120 sccm and 30-min 1100°C annealing) is shown in Fig. 6, which supports the existence of nc-Si. The lattice distance between two (111)-oriented planes of 0.63 nm observed from a Si substrate is employed as a standard ruler in this picture, which helps to estimate the smaller plane-to-plane distance of nc-Si of about 0.19 nm. According to the X-ray diffraction (XRD) data, the oriented planes of nc-Si structure are determined as (220), which is in good agreement with the reference data. The sizes of nc-Si are distributed from 4.4 nm to 5 nm from the TEM analysis.

The TRPL results of SiOₓ samples at 1100°C for different annealing time are shown as Fig. 7 and in the Table 1. A stretched exponential function: 
\[ I(t) = I_0 \exp(-t/\tau) \] was used to fit the data, in which \( \tau \) is an effective decay time of nc-Si. The luminescent lifetime lengthens from 20 to 52 µs as the nc-Si size extends from 4.0 to 4.2 nm; the nc-Si lifetime increases smoothly with the increment of nc-Si size, which has the same temper by Garcia et al.¹⁵. 

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Fig. 5. PL intensity and wavelength of nc-Si in SiOₓ film under different N₂O fluence at optimal annealing times.

Fig. 6. The TEM micrograph of the cross section of the SiOₓ film annealed at 1100°C for 30 min.
The theoretical carrier-transition equation can be simplified as \[ I_{PL} \propto \sigma \phi(t) \frac{1}{\tau_{PL}} N_{nc-Si} \] \(^{15, 19}\), and the variation of the nc-Si density can be estimated by the PL intensity \((I_{PL})\) and lifetime \((\tau_{PL})\) of nc-Si, where \(\sigma\) and \(\phi(t)\) are the emission (absorption) cross-section of nc-Si and the pumping photon flux density obtained from the pumping power, respectively, and their multiplication for different annealing-time samples are a constant. As the annealing time lengthens from 30 to 120 min, the density of nc-Si is decreased from \(8.3 \times 10^{22}\) cm\(^{-3}\) to \(1.2 \times 10^{22}\) cm\(^{-3}\), which are correlated with the evolution of measured PL and the results also reported by Augustine et al.\(^{19}\).

### 4. CONCLUSIONS

The optimized \(N_2O\) fluence for PECVD-grown SiO\(_x\) films with buried nc-Si is demonstrated. Strong PL at 550-870 nm that attributed to nc-Si has been observed in SiO\(_x\) thin films with \(N_2O\) fluence varying from 105 to 130 sccm. After annealing from 15 to 180 min, a 22-nm-redshift of the PL has been detected. The maximum PL intensity at 761 nm emitting is observed for the 30-min annealed SiO\(_x\) growing at 120-sccm \(N_2O\)-fluence. Larger \(N_2O\) fluence and longer annealing time cause a PL blueshift by 65 nm and 20 nm, respectively. Such a blueshift is attributed to shrinkage in the size of the nc-Si under the participation of dissolved oxygen atoms from \(N_2O\). We demonstrated that nc-Si need some time to grow up, and then have the maximum PL intensity and size of nc-Si, finally continuative decreased due to the participant of the O atoms from \(N_2O\) fluence. It means the nc-Si needs the different optimal annealing time to add up the maximum PL intensity under different \(N_2O\) fluences, i.e. different recipe has different annealing time even at the same annealing temperature. The (220)-oriented nc-Si with diameter ranging from 4.4 to 5.0 nm is determined. The luminescent lifetime lengthens from 20 \(\mu\)s to 52 \(\mu\)s as nc-Si size extends from 4.0 to 4.2 nm and the correlated density decreased from \(8.3 \times 10^{22}\) cm\(^{-3}\) to \(1.2 \times 10^{22}\) cm\(^{-3}\).

**Table 1.** The wavelength, size, lifetime and estimated density of nc-Si after annealing for 30, 60 and 120 min.

<table>
<thead>
<tr>
<th></th>
<th>30 min</th>
<th>60 min</th>
<th>120 min</th>
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<tbody>
<tr>
<td>Wavelength (nm)</td>
<td>761</td>
<td>751</td>
<td>743</td>
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<tr>
<td>nc-Si Size (nm)</td>
<td>4.2</td>
<td>4.1</td>
<td>4.0</td>
</tr>
<tr>
<td>Lifetime ((\mu)s)</td>
<td>52</td>
<td>31</td>
<td>20</td>
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<tr>
<td>Estimated Density (cm(^{-3}))</td>
<td>(8.3 \times 10^{22})</td>
<td>(2.4 \times 10^{22})</td>
<td>(1.2 \times 10^{22})</td>
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