10 Gbit/s tunable dual-wavelength nonreturn-to-zero–to–return-to-zero data-format transformer based on a non-direct-current-biased Fabry–Perot laser diode

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We demonstrate the conversion of a nonreturn-to-zero (NRZ)-formatted electrical data stream into a wavelength-tunable return-to-zero (RZ)-formatted optical pulse code by externally seeding a synchronously sinusoidally modulated Fabry–Perot laser diode (FPLD) with optical pseudorandom binary-sequence data at 10 Gbits/s (Gbps). The FPLD without a dc-biased current was modulated by use of a power-amplified sinusoidal wave signal (~25.6 dBm) as an NRZ-to-RZ data-format transformer, which is regeneratively amplified by a closed-loop erbium-doped fiber amplifier. The gain switching and on–off keying operation of the FPLD is initiated under the seeding of self-feedback and external-injection signals. A maximum wavelength tuning range of 30 nm with a side-mode suppression ratio of greater than 36 dB is obtained. The power penalty of the NRZ-to-RZ data-format conversion at 10 Gbps is 1.5 dB. © 2005 Optical Society of America

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1. Introduction
Versatile all-optical-controlled signal-processing techniques, such as clock frequency division by period doubling,1–4 wavelength conversion by all-optical modulation,5 upstream traffic transmission by injection locking,6 and non-return-to-zero (NRZ)–to–return-to-zero (RZ) data-format transformation,7 have been demonstrated with Fabry–Perot laser diodes (FPLDs). Owing to the considerations of electronic bandwidth and timing tolerance, the data-format converters8 are potentially required for use in interconnections in optical communication networks that employ different data formats. NRZ-to-RZ9 and RZ-to-NRZ10 data-format transformations have been demonstrated with nonlinear optical loop mirrors.11 Later, all-optical data-format conversions by use of a semiconductor optical amplifier12 and gain competition13 emerged. Moreover, a wavelength-converted and logically inverted NRZ or RZ format has also been demonstrated.7 This paper applies gain competition and feedback injection techniques to demonstrate the generation of a 10 Gbits/s (Gbps) NRZ-to-RZ data-format-transformed data pulse train from a non-dc-biased FPLD. One can generate this data-format conversion by gain switching the non-dc-biased and purely sinusoidal-wave-modulated FPLD with external optical pseudorandom binary-sequence (PRBS) seeding. The proposed FPLD-based data-format converter can implement versatile functions, including transformation of an electrical NRZ signal into an optical RZ signal, generation of a wavelength-converted and tunable RZ signal, and dual-wavelength and PRBS-encoded gain switching and on–off keying.

2. Operating Principle
In the case here, the dc-bias current of the FPLD is set to 0 mA, and the peak amplitude of the rf sinusoidal wave signal is set to just below the threshold condition in single-mode injection locking and gain switching of the FPLD. One can reduce the effective threshold current by increasing the intensity of external injection, increasing the stimulated recombina-
nation coefficient, and causing FPLD gain switching. The proposed FPLD NRZ-to-RZ data converter for performing dual-wavelength operation and wavelength conversion (Fig. 1) is based on dual-wavelength injection locking. Owing to the redshift of the gain profile under injection, the self-feedback seeding wavelength $\lambda_e$ is adjusted slightly larger than the maximal transmission intensity of the mode spectrum of the FPLD, as shown in Figs. 1(a) and 1(b). When the FPLD is externally injection locked with another wavelength $\lambda_e$, both $\lambda_f$ and $\lambda_e$ can be lased simultaneously, achieving dual-wavelength operation [Fig. 1(c)]. For achieving wavelength conversion, the wavelength deviation between $\lambda_e$ and the wavelength at the local maximal transmission intensity must be higher than that between $\lambda_f$ and the wavelength at the local maximal transmission intensity. Wavelength $\lambda_e$ is exactly coincident with one of the FPLD’s longitudinal modes within the locking range, greatly reducing the threshold current for lasing at $\lambda_e$. This situation also helps to eliminate the original $\lambda_f$ because the gain profile of the FPLD redshifts further. Eventually, the $\lambda_f$ output can be fully modulated in a manner that complements the PRBS-encoded signal at $\lambda_e$.

3. Experiment

Figure 2 shows the configuration of the FPLD-based NRZ-to-RZ data-format transformer, which is based on a hybrid injection-locked FPLD–erbium-doped fiber amplifier (EDFA) link. A commercial 1.55-μm FPLD operated at 27.5 °C with a threshold current, a mode linewidth, a central wavelength, and a mode spacing of 13 mA, 2.1 GHz, 1553.380 nm, and 1.14 nm, respectively, is employed. The FPLD was sinusoidally modulated from 1 GHz to 10 GHz with a 25.6 dBm sinusoidal signal. However, the dc-biased current of the FPLD was set to 0 mA. In this case the FPLD was neither lasing nor gain switching because the overall driving current was below the threshold. The optical output of the FPLD was fed through an 80% fiber coupler (OC1) connected to another 50–50 fiber coupler (OC2). One coupling branch of OC2 was fed through a circulator into a closed-loop EDFA, which forms an external cavity to achieve regeneratively amplified self-feedback. In the external seeding branch, the tunable laser that was externally modulated with an electrical PRBS data stream via a Mach–Zehnder intensity modulator.
modulator (M Zimmerman) was used as the external-injection seeding source for the FPLD by use of another coupling branch of OC2. In the self-feedback seeding branch, the amplified spontaneous emission from a commercial EDFA with a gain of 21 dB was filtered with an intracavity tunable narrow bandpass filter (TF), which then feeds back into the FPLD. An optical attenuator was used to control the amount of feedback from the EDFA, which optimizes the spectral and temporal characteristics of the FPLD. Two optical isolators inside the closed-loop EDFA were used to ensure unidirectional operation. The EDFA not only offers the amplified spontaneous emission to initiate the pulsation of the FPLD but also helps to amplify the mutually injection-locked pulse and the externally injection-locked pulse from the FPLD. Prior to the feedback seeding, a polarization controller was applied to improve the side-mode suppression ratio (SMSR) and ensure that the feedback polarization was coincident with the preferred TE lasing mode of the FPLD. The seeding powers used for self-feedback and external injection were approximately 540 μW in the experiment conducted here.

4. Results and Discussion

When the wavelength is properly set, the filtered amplified spontaneous emission signal of the EDFA in the self-feedback branch and the externally injected single-mode tunable laser signal concurrently help start the dual-wavelength gain switching in the purely sinusoidal-modulated FPLD (with a 0-mA dc-biased current). Gain switching can start only at the modulation peak of the FPLD because the optical feedback is relatively weak. Nevertheless, the amplified self-feedback seeding or external injection drastically decreased the effective threshold current, facilitating the suppression of spontaneous noise of the gain-switched FPLD. The FPLD cannot be gain switched when the injection wavelength is beyond the locking range (approximately 0.2 nm). With either injection, the maximum difference between the wavelengths of the dual lasing modes can be tuned to as much as 30 nm; an optimized SMSR of larger than 36 dB remains constant within a tuning range of ~14.82 nm (13 FPLD modes). A broader FPLD gain spectrum will increase the tuning range of the wavelength. Notably, the minimum tuning range is determined by the longitudinal mode spacing of the FPLD. Changing the frequency of repetition and the dc-bias conditions on the FPLD can also optimize the SMSR of the output signal.

Figure 3 reveals the corresponding pulse shapes and spectra of the FPLD in the dual-wavelength scheme. When \( \lambda_e \) is switched off [Fig. 3(b)], the peak power of \( \lambda_f \) increases 6.33 dB, as compared with that in Fig. 3(a). In Fig. 3(c), under dual-wavelength operation, the pulse widths \( \lambda_f \) and \( \lambda_e \) are 35.7 and 44.7 ps, respectively. After \( \lambda_e \) is switched off, an increase of the pulse peak power between the upper traces of Figs. 3(c) and 3(d) is the same as that observed between Figs. 3(a) and 3(b). Obtained from the self-feedback branch are a pulse width shorter and a peak power larger (by a factor of 1.2) than those obtained from the external-injection-seeding branch, owing to the regenerative amplification in the self-feedback loop and the shifted gain profile of the FPLD; Bouchoule et al. observed a similar result. Furthermore, when the output power of the rf signal increased from 25.6 to 26.1 dBm, the SMSR decreased from 36 to 32 dB because the intensity of the other side modes of the FPLD are enhanced. However, such an SMSR still exceeds the values of all similar configurations reported elsewhere. The wavelength-conversion mechanism is initiated as the wavelength of the self-feedback signal is detuned toward a shorter wavelength of the locking range. Figure 4 shows the pulse shapes and spectra of gain-switched FPLD pulses obtained with the wavelength-conversion scheme. As the \( \lambda_e \) signal is switched on [Fig. 4(a)], the gain at \( \lambda_f \) is almost depleted, and the gain-switched pulses at \( \lambda_f \) are switched off, as shown in Fig. 4(c).

![Fig. 3. Lasing spectra (upper figures) and pulse trains (lower figures) of the FPLD-based data-format transformer with the dual-wavelength scheme. (a) and (c) \( \lambda_e \) is switched on. (b) and (d) \( \lambda_e \) is switched off.](image)

![Fig. 4. Lasing spectra (upper figures) and pulse trains (lower figures) of the FPLD-based data-format transformer with the wavelength-conversion scheme. (a) and (b) The optical spectra; (c) and (d) the corresponding pulse trains when \( \lambda_e \) is on and off, respectively.](image)
However, when $\lambda_e$ is switched off [Fig. 4(b)], the system returns to single-mode ($\lambda_e$) operation, and the gain-switched pulses at $\lambda_e$ are switched on, as shown in Fig. 4(d). The respective pulse widths are 42.4 ps [Fig. 4(c)] and 31.2 ps [Fig. 4(d)]. The higher peak power of the FPLD pulse at $\lambda_e$ is due to the regenerative amplification effect in the self-feedback configuration.

In an application, one can generate a wavelength-converted, RZ-formatted optical data stream by seeding the synchronously modulated FPLD with an NRZ-format optical PRBS data signal. An electrical NRZ signal can be readily translated into an optical RZ-format signal at $\lambda_e$ or $\lambda_f$, where the signal $\lambda_e$ is out of phase with the signal $\lambda_f$, owing to the wavelength-conversion operation. Notably, signals $\lambda_e$ and $\lambda_f$ can also be in phase under a dual-wavelength operation. Such an NRZ-to-RZ data-format conversion can be applied directly to the time-division multiplexing system to increase the bit rate. For instance, Fig. 5 shows an 8-bit NRZ-formatted signal (01010101) that modulates the tunable laser and the RZ-formatted output signal of the FPLD at $\lambda_e$ with a bit rate of 1 Gbps. The highest-ever-achieved bit rate of the FPLD-based NRZ-to-RZ data-format transformer is 10 Gbps. (See inset of Fig. 5.) The bit-error-rate measurement of the synchronously modulated FPLD under the injection of the 10-Gbps NRZ data stream ($2^{23} - 1$ PRBS) is carried out, as shown in Fig. 6. The converted RZ is detected by a lightwave clock–data receiver, incorporating a high-gain avalanche photodiode, a gain-controlled amplifier, and a clock–data recovery circuit. The received power is measured in the back-to-back NRZ data, and the back-to-back NRZ data are converted into RZ data at 10 Gbps. The power penalty is measured in the back-to-back NRZ data, and the back-to-back NRZ data are converted into RZ data at 10 Gbps. The power penalty at the $10^{-9}$ bit-error-rate level is 1.5 dB. Notably, an additional dc-bias current is required to further extend the operating bandwidth of the FPLD.

EDFA in the external cavity here yields a larger SMSR (>36 dB); (2) the FPLD-based data-format transformer is operated in the pulsed mode, by one’s setting the FPLD under a non-dc-current driving condition; and (3) a higher wavelength switching speed of as much as 10 GHz is reported, which compares favorably with the 300 KHz of the previous system. Nonetheless, the carrier and photon dynamics of the FPLD continue to limit the maximum speed of wavelength conversion.

5. Conclusions
A novel FPLD-based configuration for translating an NRZ-format electrical signal into a wavelength-tunable, RZ-format optical signal was demonstrated. Gain switching the FPLD with a non-dc current and combining external-injection and self-feedback seeding promotes versatile functionality, including the dual-wavelength operation, wavelength-conversion operation, and NRZ-to-RZ data-format transformation. A Mach–Zehnder intensity modulator is used to encode the injection light by use of an NRZ-formatted PRBS pattern, yielding a 10-Gbps RZ-format optical signal from the proposed FPLD-based data transformer. Properly setting the power and wavelength of the injection signal allows dual-wavelength and wavelength-conversion operations to be demonstrated. A maximum tuning range and SMSR of 30 nm and 36 dB are observed, respectively. The power penalty of 1.5 dB at 10 Gbps is measured for the NRZ-to-RZ data-format conversion. This configuration overcomes the drawbacks of the previous setups, such as the wavelength-dependent path in the grating-based self-feedback system, the difficulty in selecting identical FPLDs for an external-injection system, and the need for alignment of the optical path in the grating-based configuration.

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References


